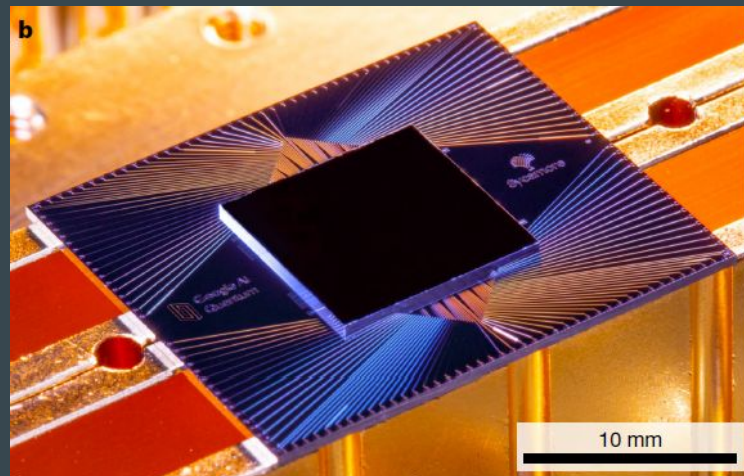


Technical challenges and realizations of quantum qubits: the example of superconducting charge qubits

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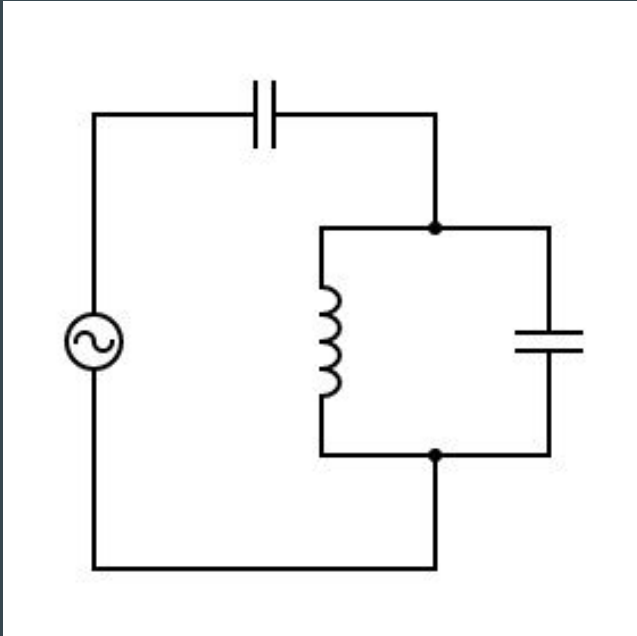
The Sycamore processor [F. Arute et al, 2019]

The challenges

- Scalability
- Protected from decoherence (large relaxation time T_1 and dephasing time T_2)
- Protected from noise
- Implementation of single-qubit and two-qubit quantum gates
- Measurement apparatus

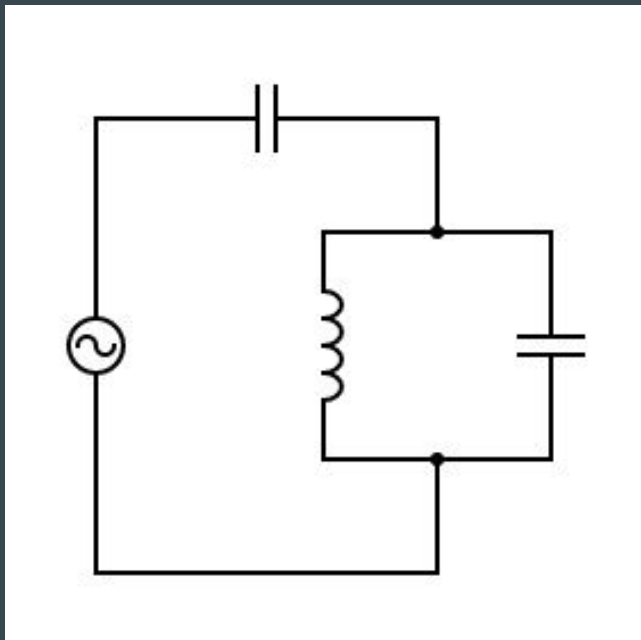
Superconducting qubits: the Cooper Pair Box (CPB)

$$\hat{H} = Q^2 / 2C + \Phi^2 / 2L$$

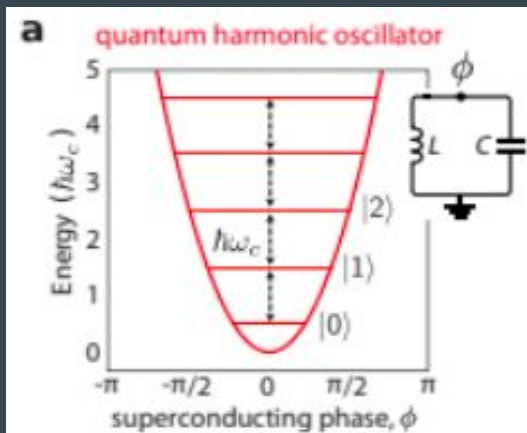


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[D.Bernal et al. , 2020]

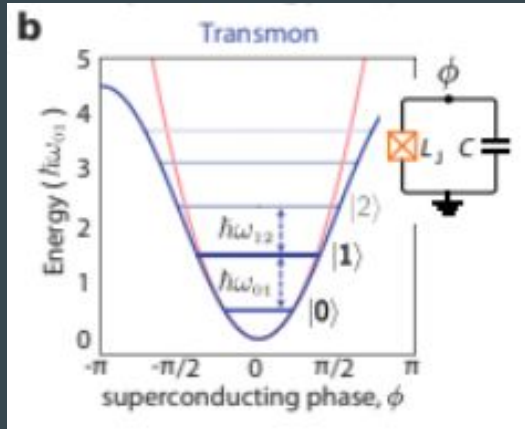
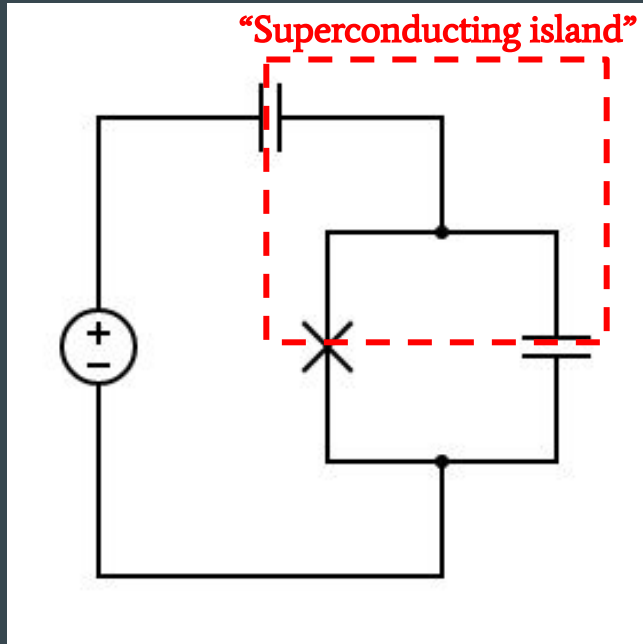


Superconducting qubits: the Cooper Pair Box (CPB)

$$\hat{H} = E_C (\hat{n} - n_g)^2 - E_J \cos(\hat{\varphi})$$

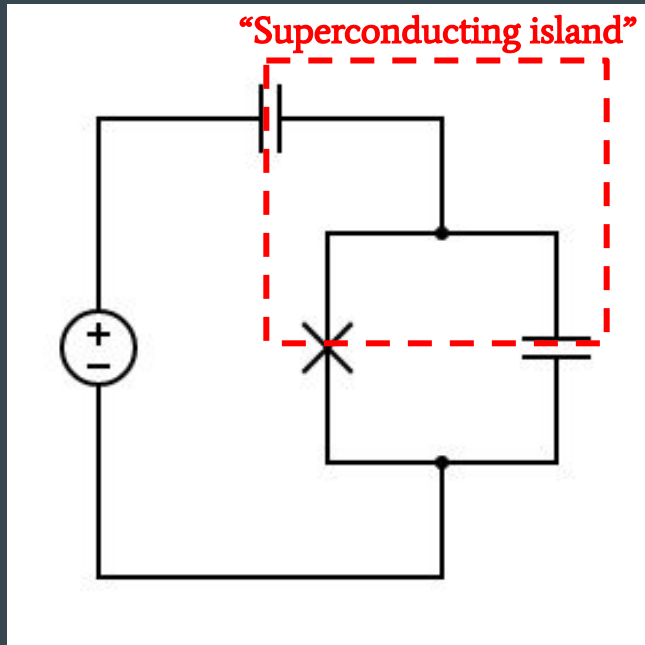
$$\hat{Q} = (2e)^2 (\hat{n} - n_g)^2$$

[D.Bernal et al., 2020]



Superconducting qubits: the Cooper Pair Box (CPB)

$$\hat{H} = -E_J \sum_n \left(|n+1\rangle\langle n| + |n\rangle\langle n+1| \right) + E_C \sum_n (n - n_g)^2 |n\rangle\langle n|$$



$$E_C = (2e)^2 / 2C_\Sigma$$

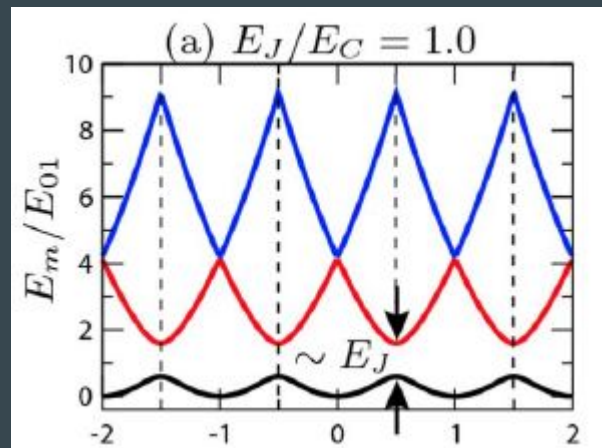
$$n_g = C_g V_g / 2e$$

Superconducting qubits: the Cooper Pair Box (CPB)

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To obtain a qubit:
tuning of E_c and n_g to
maximize
anharmonicity

$$\alpha = \Delta E_{12} - \Delta E_{01}$$



[J.Koch et al., 2007]

$$E_J \sim \text{constant}$$

$$E_C = (2e)^2 / 2C_\Sigma$$

$$n_g = C_g V_g / 2e$$

Superconducting qubits: the Cooper Pair Box (CPB)

$$\hat{H} = -E_J \left(|1\rangle\langle 0| + |0\rangle\langle 1| \right) \\ + E_C \left(n_g^2 |0\rangle\langle 0| - (1 - n_g)^2 |1\rangle\langle 1| \right)$$

Superconducting qubits: the Cooper Pair Box (CPB)

$$\hat{H} = K(n_g)\mathbb{I} - E_c(n_g - 1/2)\sigma_z + E_J\sigma_x/2$$

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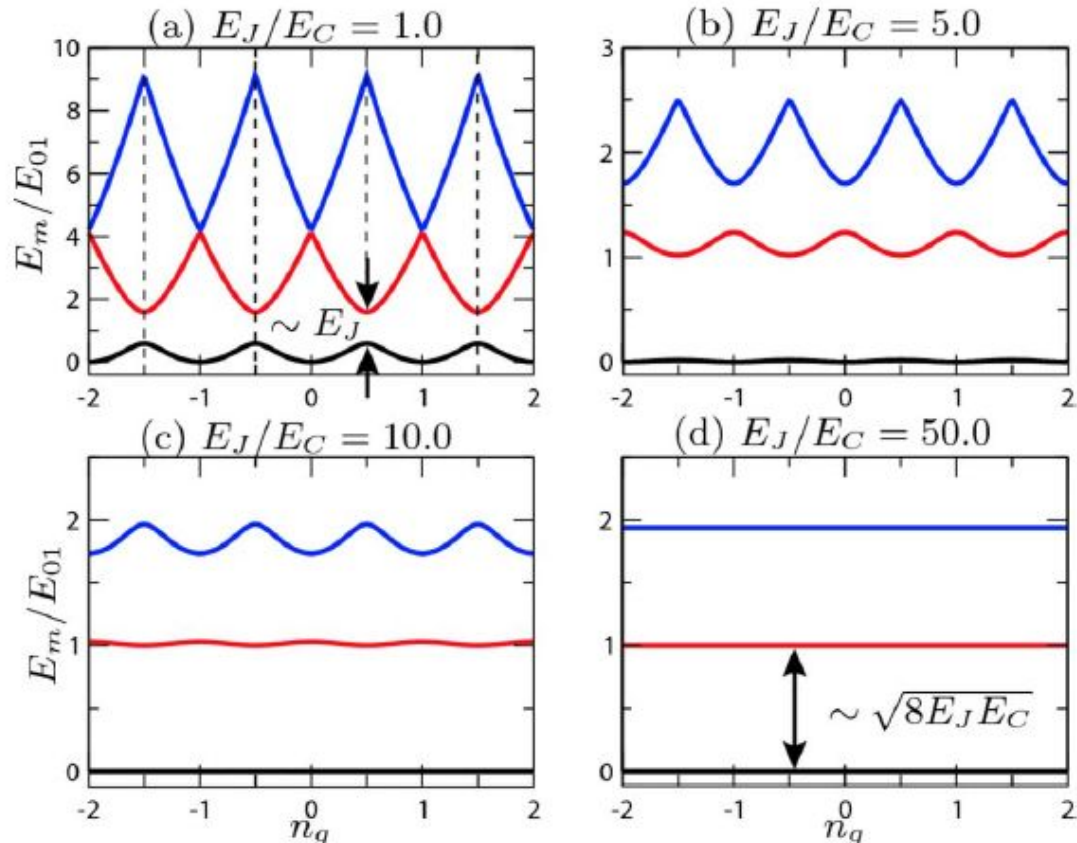
n_g can be used as a control to change the nature of H , henceforth implementing single-qubit quantum gates:

- adiabatic change (e.g. from $n_g=0$ to $n_g=1/2$ to obtain $|+\rangle$)
- fast change of n_g leads to Rabi-like dynamics; then we control the onset time t to obtain the desired state.

Superconducting qubits: the Cooper Pair Box (CPB)

Problem: sensitivity to charge noise, affecting the value of n_g , inducing small dephasing time T_2 .

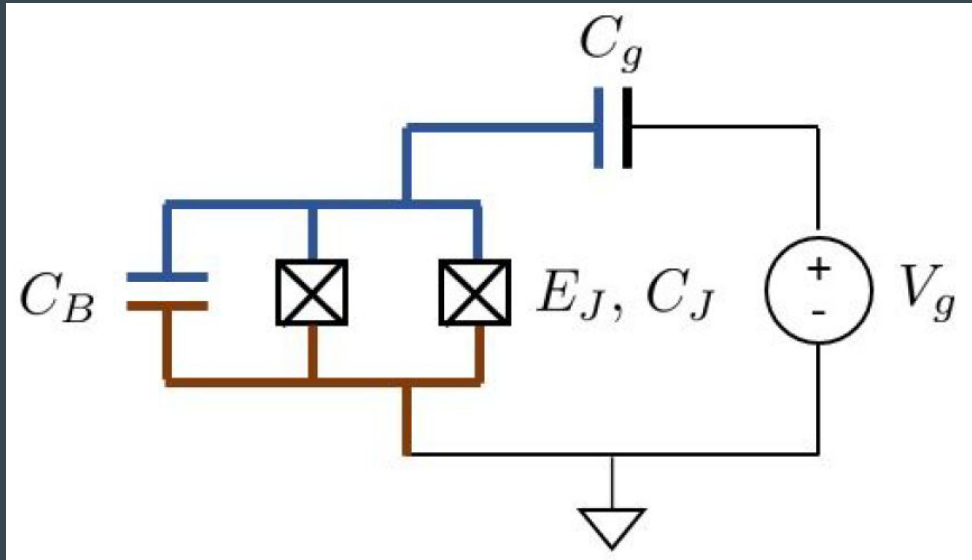
Superconducting qubits: the Transmon



In the transmon regime, E_J/E_C is tuned so that $E(n_g)$ dispersion becomes negligible, and some anharmonicity is preserved. n_g becomes irrelevant.

[J.Koch et al. , 2007]

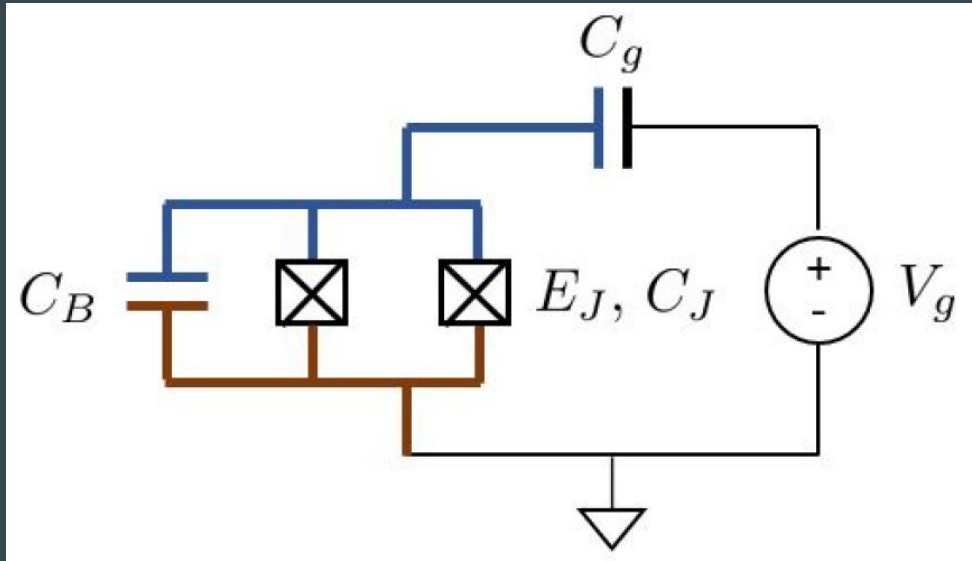
Superconducting qubits: the Transmon



In practice, E_C is tuned by increasing total capacitance.

[T. Roth et al. , 2021]

Superconducting qubits: the Transmon



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In practice, E_C is tuned by increasing total capacitance.

New problem is the weak anharmonicity: risks of leakage. Pulse shaping allows fast transitions with fine frequency control.

Superconducting qubits: the Transmon

The transmon qubits can be coupled (entangled), through a resonator, which is expressed by a Jaynes-Cummings Hamiltonian:

$$\hat{H} = \hat{H}_r + \sum_{j=1,2} \hat{H}_j + \sum_{j=1,2} g(a^\dagger \sigma_j^- + a \sigma_j^+)$$

Superconducting qubits: the Transmon

The transmon coupled to a circuit resonator can be expressed, after some manipulations, by a (approximated) Jaynes-Cummings Hamiltonian:

$$\hat{H} = (\omega_r - \chi|1\rangle\langle 1| + \chi|0\rangle\langle 0|)a^\dagger a + (\omega_1 + \chi)\sigma_z$$

Superconducting qubits: the Transmon

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The resonator frequency depends on the transmon state; it is then measured by microwave spectroscopy.

Superconducting qubits: advantages

- Easy to embed into a larger circuit
- 1D junctions effectively smaller than 3D cavity
- Easy control transition frequencies: protection from thermal noise
- Macroscopic quantum state with high coupling strength
- Straightforward measurement

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Thank you.