

Peculiarities and applications of narrow-line magneto optical traps for neutral atoms

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Contents

Generalities

The doppler cooling and trapping

Narrow line cooling

Experiments

History

- 1925 Bose and Einstein predicted a new state of matter
- 1975 Hänsch and Schawlow and Wineland and Dehmelt proposed the optical molasses
- 1982 Phillips and Metcalf realized the Zeeman slower
- 1989 Chu realized the first magneto optical trap;
- 1997 Chu, Cohen-Tannoudji and Phillips received the Nobel Prize;
- 2001 Cornell, Wieman, and Ketterle received the Nobel Prize

Why are MOT cool

Mot characteristics:

- Low temperatures
- high densities
- relatively simple and robust
- are the workhorse of most cold atom experiments.

Application:

- Degenerate quantum gases
Lanthanides: Er, Dy, Yb...
- Atomic clocks
bosonic alkaline earth elements ^{88}Sr ...

The scattering rate

The rate at which an atom at rest absorbs photons is given by:

$$R_s = \frac{\gamma}{2} \frac{s}{1 + s + 4\delta^2/\gamma^2}$$

where $\delta = \omega_L - \omega_A$, γ is the transition linewidth and s is the saturation parameter

- an atom absorbs a photon with $\hbar\omega_L$ energy
- emits a photon with $\hbar\omega_A$ energy
- the atom momentum changes by $\vec{p} = \hbar\vec{k}$
- the atom is subjected to a force $\vec{F}_s = \hbar\vec{k}R_s$

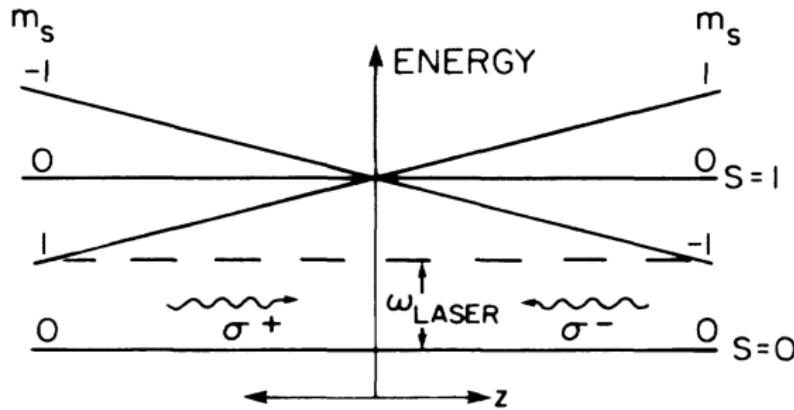
Principles of operation

The detuning δ is velocity dependent via the Doppler effect:

$$\omega'_L = \omega_L - \vec{k} \cdot \vec{v} \rightarrow \vec{F} = -\alpha \vec{v}$$

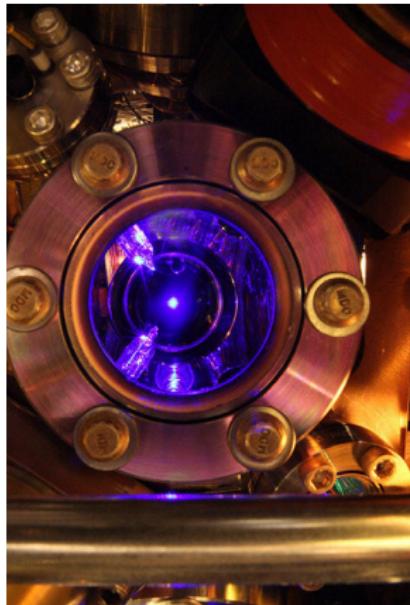
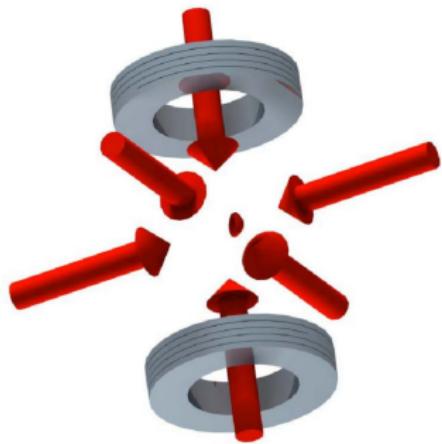
The detuning can be made position dependent using a magnetic field gradient:

$$\omega'_A = \omega_A + \frac{\mu \nabla B z}{\hbar}$$



Credit: E. L. Raab et al Phys. Rev. Lett. 59, 2631 (1987)

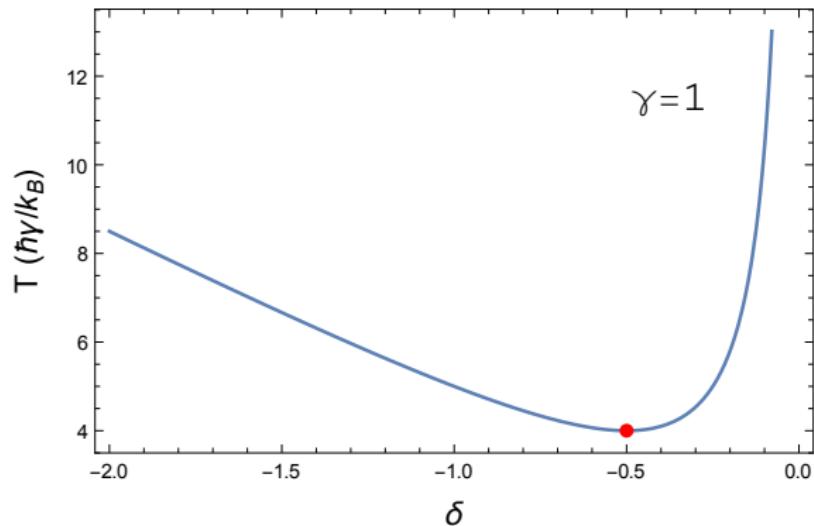
Optical Scheme



Credit: Dreon et. al. J. Phys. B: At. Mol. Opt. Phys. 50 (2017)
Credit: The Ye group and Brad Baxley, JILA

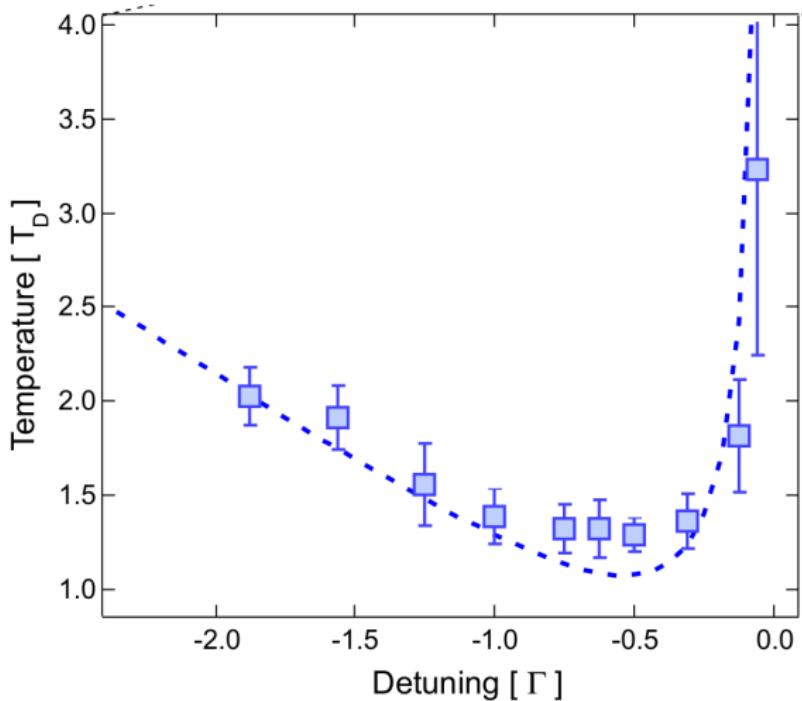
Doppler temperature

$$T = \frac{\hbar\gamma}{8k_B} \left(\frac{\gamma}{|\delta|} + 4 \frac{|\delta|}{\gamma} \right) \quad T_{\min} \propto \frac{\gamma}{2} \quad \text{for} \quad \delta = -\frac{\gamma}{2}$$



Experimental data in 3D

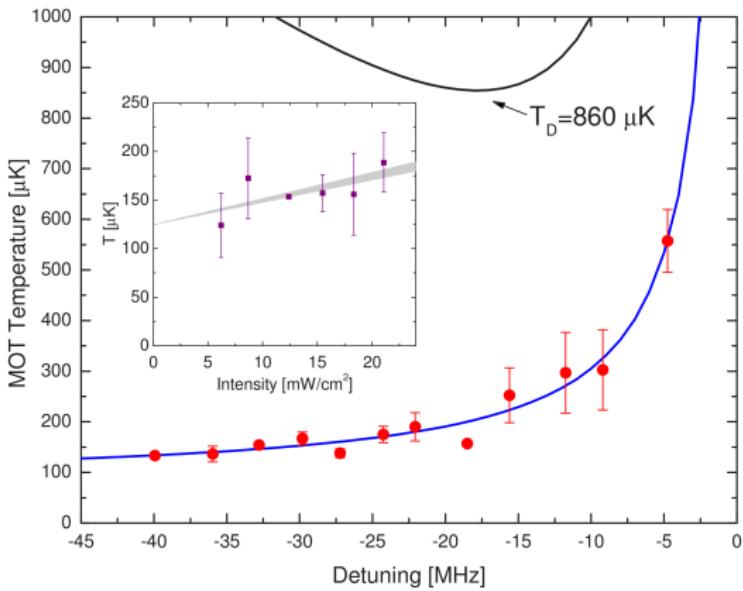
${}^4\text{He}^*$
 T_D 35.75 μK



Credit: R. Chang et. al. Phys. Rev. A 90 (6), pp.063407 (2014)

Sub-Doppler cooling of Er

- $\gamma \simeq 36 \text{ MHz}$
- $T_D = 860 \mu\text{K}$
- $T \simeq 100 \mu\text{K}$
- 10^4 atoms



Credit: A. J. Berglund et. al. Phys. Rev. A 76, 053418 (2007)

Time scales

$$E_r = \frac{\hbar^2 k^2}{2m} \quad \Rightarrow \quad \omega_r = \frac{\hbar k^2}{2m}$$

Atom internal state evolve according to γ :

$$t_{\text{int}} = 1/\gamma$$

Atom external time:

$$kv = \gamma \quad \text{and} \quad a = v_r \gamma \implies t_{\text{ext}} = 1/kv_{\text{rec}} = \hbar/2E_r$$

Broadband

$$t_{\text{ext}} \gg t_{\text{int}}$$

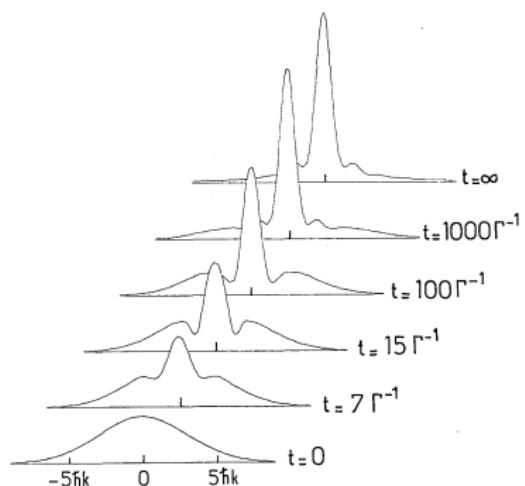
$$\gamma \gg 2\omega_r$$

Narrow line

$$t_{\text{ext}} \simeq t_{\text{int}}$$

$$\gamma \simeq 2\omega_r$$

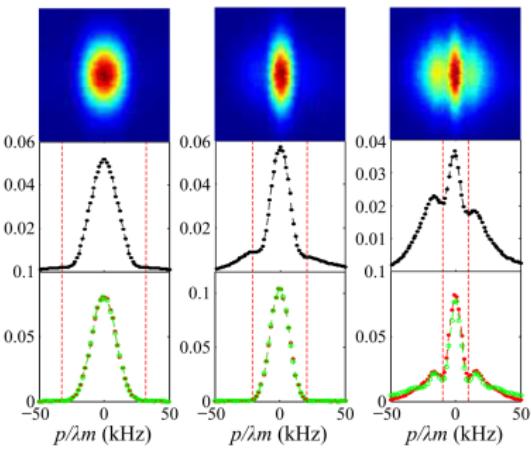
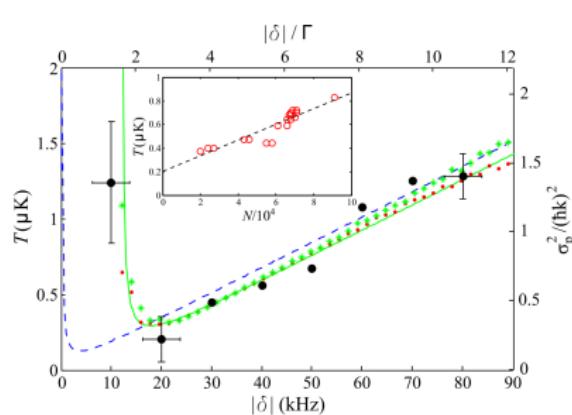
Narrow line MOT



- Optimal detuning $\delta \simeq 3\omega_r$
- Predicted final temperature $\simeq 1.5E_r/k_B$
- Semiclassical theory is no longer valid
- velocity distribution is a power law $P(v) \propto |v|^{-2\frac{|\delta|}{\omega_r}}$
- small capture velocity
- gravity can't be neglected:
$$R = \frac{\hbar\gamma k}{2} \frac{1}{mg} \simeq 1$$

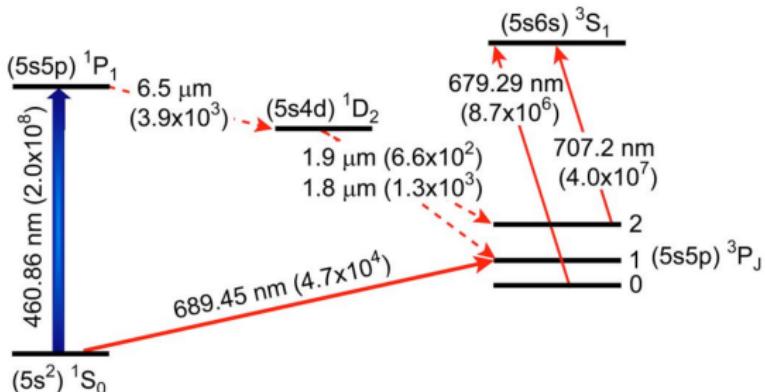
Credit: Castin et. al. J. Opt. Soc. Am. B/Vol. 6, No. 11 (1989)

Narrow line ^{88}Sr MOT



Credit: Chalony et. al. PRL 107, 243002 (2011)

Narrow line ^{88}Sr MOT



$$461 \text{ nm: } ^1S_0 \rightarrow ^1P_1$$

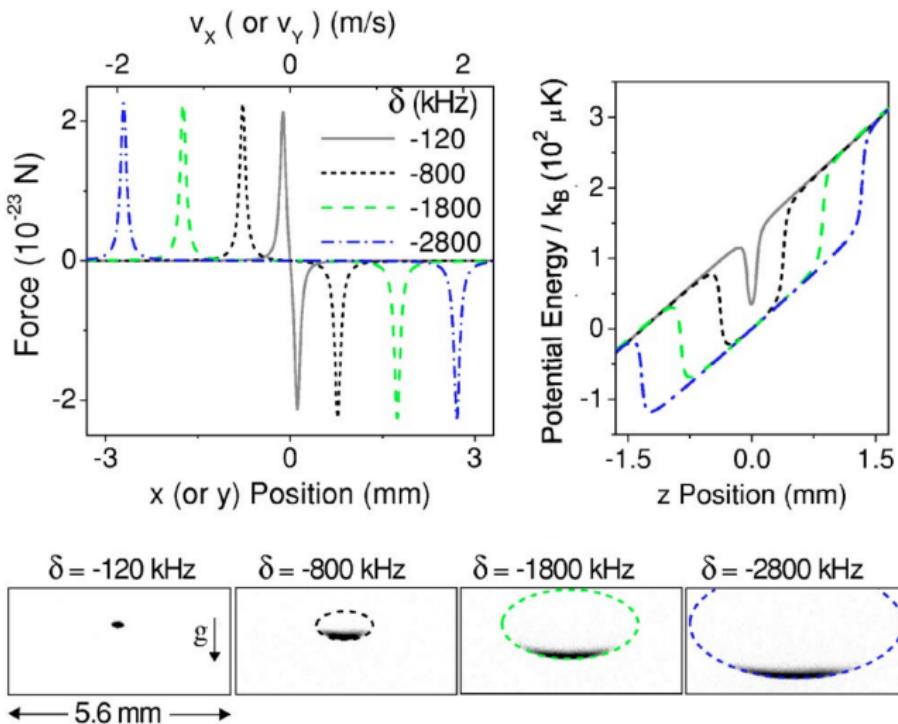
- $\gamma = 32 \text{ MHz}$
- $\omega_r/2\pi = 10.7 \text{ kHz}$
- $T_D = 768 \mu\text{K}$
- $\frac{\hbar\gamma k}{2mg} = 43 \times 10^3$

$$689 \text{ nm: } ^1S_0 \rightarrow ^3P_1$$

- $\gamma = 7.5 \text{ kHz}$
- $\omega_r/2\pi = 4.7 \text{ kHz}$
- $T_D = 0.3 \mu\text{K}$
- $\frac{\hbar\gamma k}{2mg} = 16$

Credit: Loftus et. al. Phys. Rev. Lett. 93, 073003 (2004)

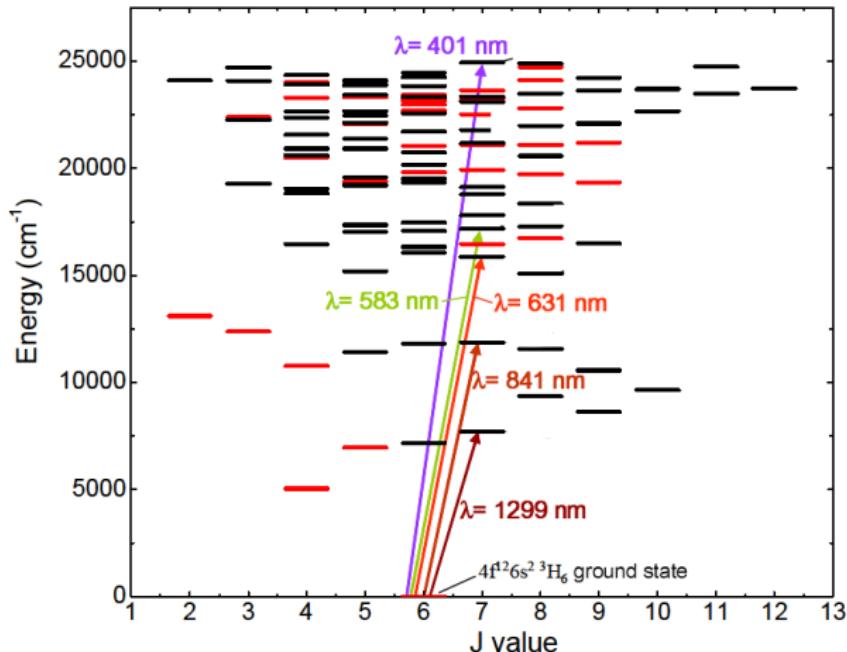
Narrow line ^{88}Sr MOT



Credit: Loftus et. al. PRA 70, 063413 (2004)

Lanthanides

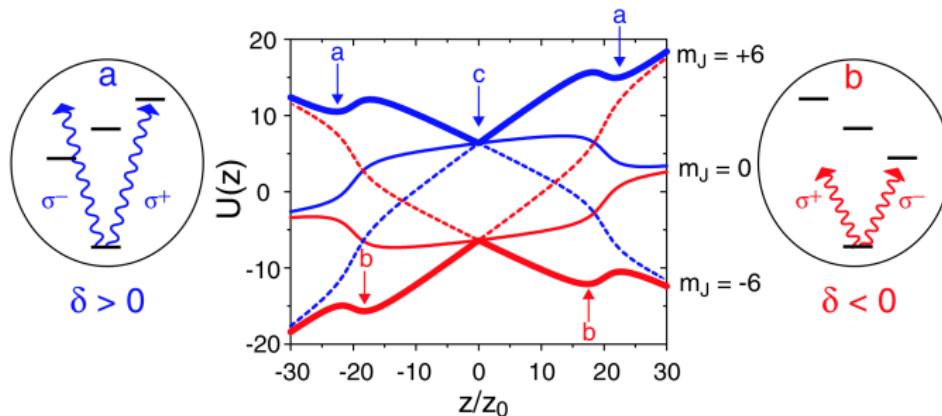
λ (nm)	401	583	841
γ (kHz)	36	190	8
T_D (μK)	660	4.8	0.19



^{166}Er MOT

$$\mu = 7\mu_B \implies F = F_s - \mu |\nabla B| - mg, \quad \mu = g m_j \mu_B$$

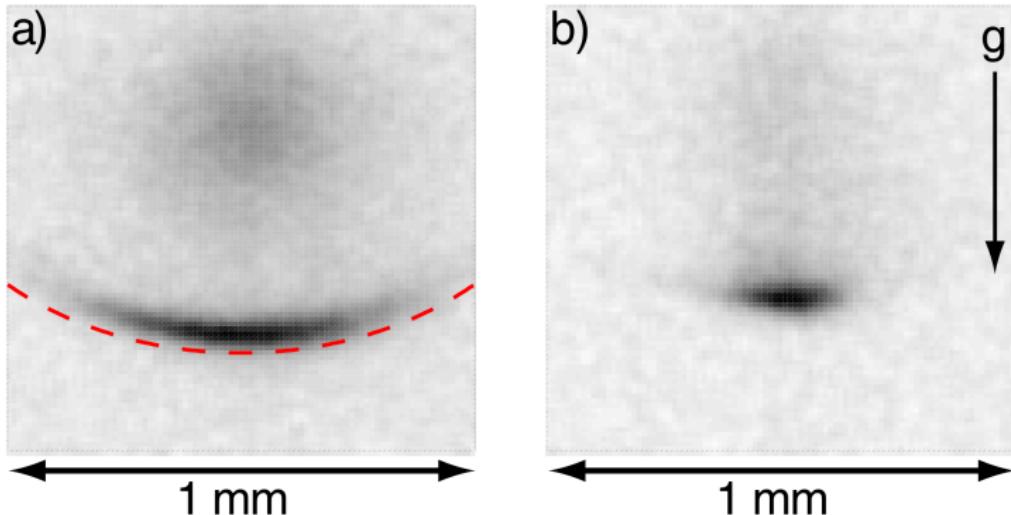
$$\omega_r / 2\pi = 1.7 \text{ kHz}, \quad \hbar \gamma k / 2mg = 7$$



- the $m_j = j$ state is a “low field seeking state”
- For $|\nabla B| \neq 0 \quad F \rightarrow F_s + g m_j |\nabla B|$

Credit: Berglund et. al. PRL 100, 113002 (2008)

^{166}Er MOT



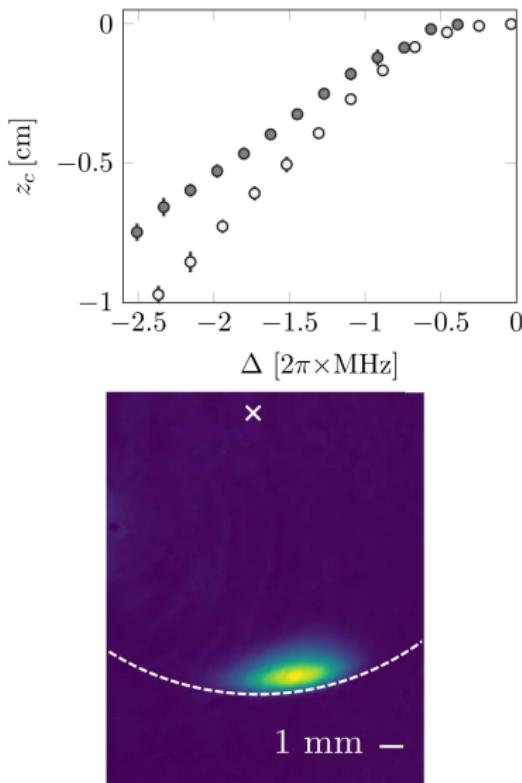
$$T = 2\mu K$$

Atoms are polarized in the $J = 6$ $m_j = 6$ state

Credit: Berglund et. al. PRL 100, 113002 (2008)

^{162}Dy MOT

- 626 nm 136 kHz transition
- $\omega_r/2\pi = 3.1\text{kHz}$, $\hbar\gamma k/2mg = 171$
- $\delta = 30\gamma$
- 3×10^8 atoms at $T = 20\mu K$
- Mot vertical position self-adjust for $F_s = mg$



Credit: Dreon et. al. J. Phys. B: At. Mol. Opt. Phys. 50 (2017)

Conclusions

Narrow line MOT:

- Provide low temperatures
- Gravity plays an important role
- Require preliminary cooling stages