Optical quantum sensors for brain monitoring and imaging

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Brain monitoring: until recently, dominated by static neuroimaging

- 3 main traditional techniques:
- \bullet electroencephalography (EEG): time resolution \sim ms; spatial res. \sim cm
- \bullet magnetoencephalography (MEG): time resolution \sim ms; spatial res. \sim mm

 \bullet functional magnetic resonance imaging (fMRI): time resolution up to \sim 0.1 ms; spatial res. ∼ mm

 \rightarrow better SNR due to blood oxygen level-dependent contrast technique, recently also from introduction of spin-relaxation free optical magnetometers

Limitation: large, expensive, patient has to stay still inside the scanner!!

Not wearable \rightarrow spatial resolution limited by distance between sensor and brain

Optical techniques: promise to be more flexible

Optical techniques: general characteristics

Sensing principle: functional changes in tissue lead to changes in static neuroimaging

Main advantages over traditional techniques:

- \bullet smaller, typically more cost-effective
- \bullet wearable \to adaptable to different environments \to better suited to clinical applications
- \bullet Can exploit different wavelenghts to monitor various parameters in parallel

Disadvantage: strong scattering and absorption in human skull and scalp

 \rightarrow typical transmission for 7-11mm-thick skull: ~2% @1064nm

 \rightarrow Some techniques only usable on soft tissues, brains of infants/small animal models

 \rightarrow difficult to use for morphological scanning

4 main techniques for label-free, non-invasive optical brain monitoring:

 functional near-infrared spectroscopy (fNIRS): monitor cortical neural activity measuring dynamics of hemoglobin, but limited spatial resolution

 \bullet diffuse correlations spectroscopy (DCS): measure blood flow thanks to scattering of photons by blood cells, but limited spatial resolution

 photoacoustic imaging (PAI): high contrast and spatial resolution, limited by thick skill effect (ultrasound absorption ~20 dB/cm)

optical coherence tomography (OCT): extremely accurate spatial information, good temporal resolution, extremely limited penetration depth

Consequence: only first two techniques usable for human brain monitoring!

Optical techniques: overview /2

 $CMRO₂$ = cerebral metabolic rate of oxygen; HbT = concentration of hemoglobin; LD = laser diode; $BFI = blood flow index$; $CBF = cerebral blood flow$; $PD = photodiode$; $APD = avalanche$ $photodiode$; $PZT = piezoelectric transducer$

functional Near-InfraRed Spectroscopy (fNIRS)

Sensing principle: wavelength-dependent absorption and scattering due to different energy level structure of functional molecules (e.g. hemoglobin)

Typically uses 2-3 wavelenghts in physiological window 650-950 nm \rightarrow skin, tissues, bones ∼ transparent

functional Near-Infrared Spectroscopy (fNIRS)

• Clinical applications:

preoperative localization of tumors/cancer therapy or neurodegenerative disease monitoring

mapping of functional areas before surgery

• Uses high-sensitivity photodiodes/photomultipliers to detect back-scattered light along "elliptical pathways"

 \rightarrow need for high-powered sources (LED, LD) near the brain

$$
\mu(z) = \sum_{i=1}^N \mu_i(z) = \sum_{i=1}^N \sigma_i n_i(z)
$$

 $\mu(z)$ attenuation coefficient at point z, σ_i attenuation cross-section of species i, $n_i(z)$ number density of species i

$$
T=e^{\int_0^d\mu(z)\,dz}=10^{-A}
$$

 T transmittance coefficient, A absorbance coefficient, I optical path length

$$
OD = log_{10}(\frac{l_0}{l_{in}}) = \mu_a \cdot [X] \cdot I \cdot DPF + G
$$

OD optical density, I_0 measured light intensity, I_{in} incident light intensity, $[X]$ cromatophore concentration, DPF differential path length factor (determined through measurement of the mean time of ight of picosecond pulse or by frequency-domain NIRS at 4 different wavelengths), G geometrical parameter for the light's scattering properties (treated as constant)

By measuring the differentials in time, we obtain information on the changes in cromatophore concentrations:

$$
\Delta[X] = \frac{\Delta OD}{\mu d}
$$

 $(d$ total corrected photon path length)

 \rightarrow measuring at at least two wavelengths λ_1 and λ_2 (one below and one above the crossing point):

$$
\begin{pmatrix} \Delta OD_{\lambda_1} \\ \Delta OD_{\lambda_2} \end{pmatrix} = \begin{bmatrix} \mu_{\lambda_1}^{Hb}d & \mu_{\lambda_1}^{HbO_2}d \\ \mu_{\lambda_2}^{Hb}d & \mu_{\lambda_2}^{HbO_2}d \end{bmatrix} \begin{pmatrix} \Delta[X]^{Hb} \\ \Delta[X]^{HbO_2} \end{pmatrix}
$$

fNIRS - techniques /1

3 different fNIRS techniques:

- Continuous Wave technique
- Frequency-Domain technique
- Time-Domain technique

Absolute concentrations obtainable from total photon path length

Continuous Wave technique

low cost, miniaturized and wearable

doesn't measure time-of-flight, so only monitors relative changes in the ratio HbO2/Hb

can't distinguish between absorption and scattering changes

amplitude of different wavelengths modulated at kHz/10s of kHz frequency (i.e. using acousto-optic modulators) to distinguish them

better SNR due to use of lock-in amplification

Lock-in amplification: If we multiply a periodic signal by a cosine/sine reference signal, the average amplitude is: $U_{out} = \frac{1}{2} V_{sign} \cdot cos \theta$

 V_{sien} measured signal, θ dephasing btw reference and signal

$$
\rightarrow \text{using two detectors: } \begin{cases} X = \frac{1}{2} V_{\text{sign}} \cdot \cos \theta \\ Y = \frac{1}{2} V_{\text{sign}} \cdot \sin \theta \end{cases} \rightarrow \begin{cases} U_{\text{out}} = \sqrt{X^2 + Y^2} \\ \theta = \arctan \frac{Y}{X} \end{cases}
$$

fNIRS - techniques /3

Improvement due to lock-in amplication: because noise is wide-bandwidth $P_{noise} = k_B \cdot T_{\text{eff}} \cdot \Delta \nu$, T_{eff} effective temperature of detection chain, $\Delta \nu$ measuring bandwidth

 \rightarrow measuring over small $\varDelta\nu$, we improve SNR... but we also limit time resolution $\sim 1/\varDelta\nu$

Frequency-Domain technique

intensity modulated CW laser, but at 10s/100s of MHz

allows to measure attenuation, phase shift and average path length

 \rightarrow allows to measure absolute values for Hb/HbO₂ due to estimation of reduced scattering coefficients

Time-Domain technique

uses very short (\sim ps) near infrared laser pulses

allows to obtain path length from time-of-flight \rightarrow direct measure of scattering and attenuation \rightarrow allows to measure absolute values for Hb/HbO₂

more costly and bigger, but becoming smaller, more affordable and robust

needs photon counters

Diffuse Correlations Spectroscopy (DCS)

Sensing principle: measure the Blood Flow Index (BFI) from optical properties of tissue (scattering from moving "speckles" i.e. blood cells)

Uses correlation diffusion equation (obtained from radiation transport equation $+$ some approx.)

 \rightarrow allows to obtain hemodynamics from measurement of temporal autocorrelation function of the electric field $G_1(\tau)$

 $G_1(\tau)$ related to fraction of moving scatterers over all scatterers α , to effective photon diffusion coefficient D_{γ} , to absorption and reduced scattering coefficients μ_a and μ_s

 \rightarrow limited by estimation of μ_s ', less sensitive to μ_s

Source: high coherence laser

Detectors: single-photon counting avalanche photodiodes + photon correlator

Clinical applications: high potential in cancer therapy monitoring (hemodynamics before/after)

Diffuse Correlations Spectroscopy (DCS)

• Electric field and intensity autocorrelation functions:

$$
A(\tau) = \int_{-\infty}^{\infty} E(t) E^*(t - \tau) dt
$$

$$
I(\tau) = \int_{-\infty}^{\infty} I(t) I(t - \tau) dt
$$

 \rightarrow measure of normalized intensity autocorrelation

$$
g_2(\tau)=\tfrac{\langle I(t)I(t-\tau)\rangle}{\langle I\rangle^2}
$$

 \rightarrow then we obtain $g_1(\tau)$ from the Siegert relation:

$$
G_2(\tau)=\langle I\rangle^2+\beta|G_1(\tau)|^2
$$

$$
\rightarrow g_1(\tau)=\sqrt{g_2(\tau)-1/\beta}
$$

 β constant dependent on number of speckles detected, coherence length and stability of the laser

 \bullet Correlation diffusion equation:

 $[D_{\gamma}\nabla^2 - v\mu_a - \frac{1}{3}v\mu_s' k_0^2 \langle \Delta r^2(\tau) \rangle] G_1(r,\tau) = -v S(r)$

 $\mu_s'=\mu_s(1-\langle cos\theta\rangle)=1/l^*$ reduced scattering coefficient $(l^*$ reduced scattering length), $\mu_{\textit{a}}=1/l_{\textit{a}}$ absorption coefficient ($l_{\textit{a}}$ absorption length), $D_{\gamma}=\textit{v}/3\mu_{\textit{s}}'$ photon diffusion coefficient, v speed of light in the medium, k_0 wave vector of light, r spatial position, $S(r)$ light source distribution

 $\langle\varDelta r^2(\tau)\rangle=\langle\varDelta V^2\rangle\tau^2$ for random flow with mean square velocity $\langle\varDelta V^2\rangle$

We can calculate the approximate impulse response: for semi-infinite boundary conditions, the Green function is

$$
G_1(r,\tau) = \frac{3\mu_s'}{4\pi} \left[\frac{e^{-K(\tau)\tau_1}}{r_1} - \left[\frac{e^{-K(\tau)\tau_b}}{r_b} \right] \right]
$$

$$
r_1 = \sqrt{1/\mu_s'^2 + r^2}; \ r_b = \sqrt{(2z_b + 1/\mu_s')^2 + 1/\mu_s'^2 + r^2}; \ K(\tau) = \sqrt{3\mu_a\mu_s' + 6\mu_s^2 k_0^2 \alpha \tau BFI}
$$

Solve by fitting experimental points for μ_s' , μ_a and BFI

 \rightarrow Additional constraint helpful! measure of $\mu_\text{\emph{eff}}=\sqrt{3\mu_a\mu_s'}$ by pulsed fNIRS

DCS: iterative fitting

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fNIRS: finger tapping demonstration

One sensor positioned on contralateral sensorimotor cortex, one on area unrelated to movement planning and execution; subject told to finger tap between vertical bars in graphs

DCS: models and measures of $g_1(\tau)$

Measure of blood flow in rat when flow ensured by heart/by heart pump; same measure for rat alive/dead

Basic experimental schemes of PAI and OCT

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Thank you for your attention!!