



UNIVERSITÀ  
DI SIENA  
1240

# Statistical analysis to set upper limits on Dark Matter (DM) parameters

Abhishek

Università degli Studi di Siena





# Search for DM

- **Cosmological  $\Lambda$ CDM model**
- Weakly Interacting Massive Particles (WIMPs)
- Particles are stable, massive, and interacts only through weak and gravitational interactions.
- Different approaches: Collider search, Direct detection and Indirect detection.
- Indirect: WIMPs annihilate into SM particles (bosons, quarks, leptons)



# Search for dark matter

- **Cosmological  $\Lambda$ CDM model**
- **Weakly Interacting Massive Particles (WIMPs)**
- Particles are stable, massive, and interacts only through weak and gravitational interactions.
- Different approaches: Collider search, Direct detection and Indirect detection.
- Indirect: WIMPs annihilate into SM particles (bosons, quarks, leptons)



# Search for dark matter

- **Cosmological  $\Lambda$ CDM model**
- **Weakly Interacting Massive Particles (WIMPs)**
- **Particles are stable, massive, and interacts only through weak and gravitational interactions.**
- Different approaches: Collider search, Direct detection and Indirect detection.
- Indirect: WIMPs annihilate into SM particles (bosons, quarks, leptons)



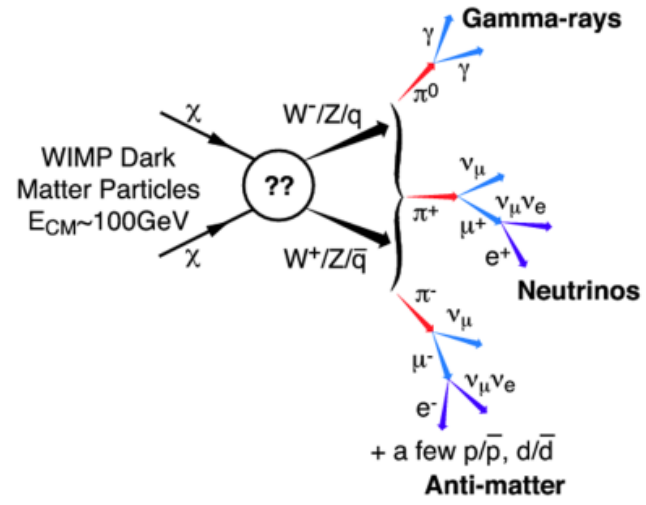
# Search for dark matter

- **Cosmological  $\Lambda$ CDM model**
- **Weakly Interacting Massive Particles (WIMPs)**
- **Particles are stable, massive, and interacts only through weak and gravitational interactions.**
- **Different approaches: Collider search, Direct detection and Indirect detection.**
- Indirect: WIMPs annihilate into SM particles (bosons, quarks, leptons)

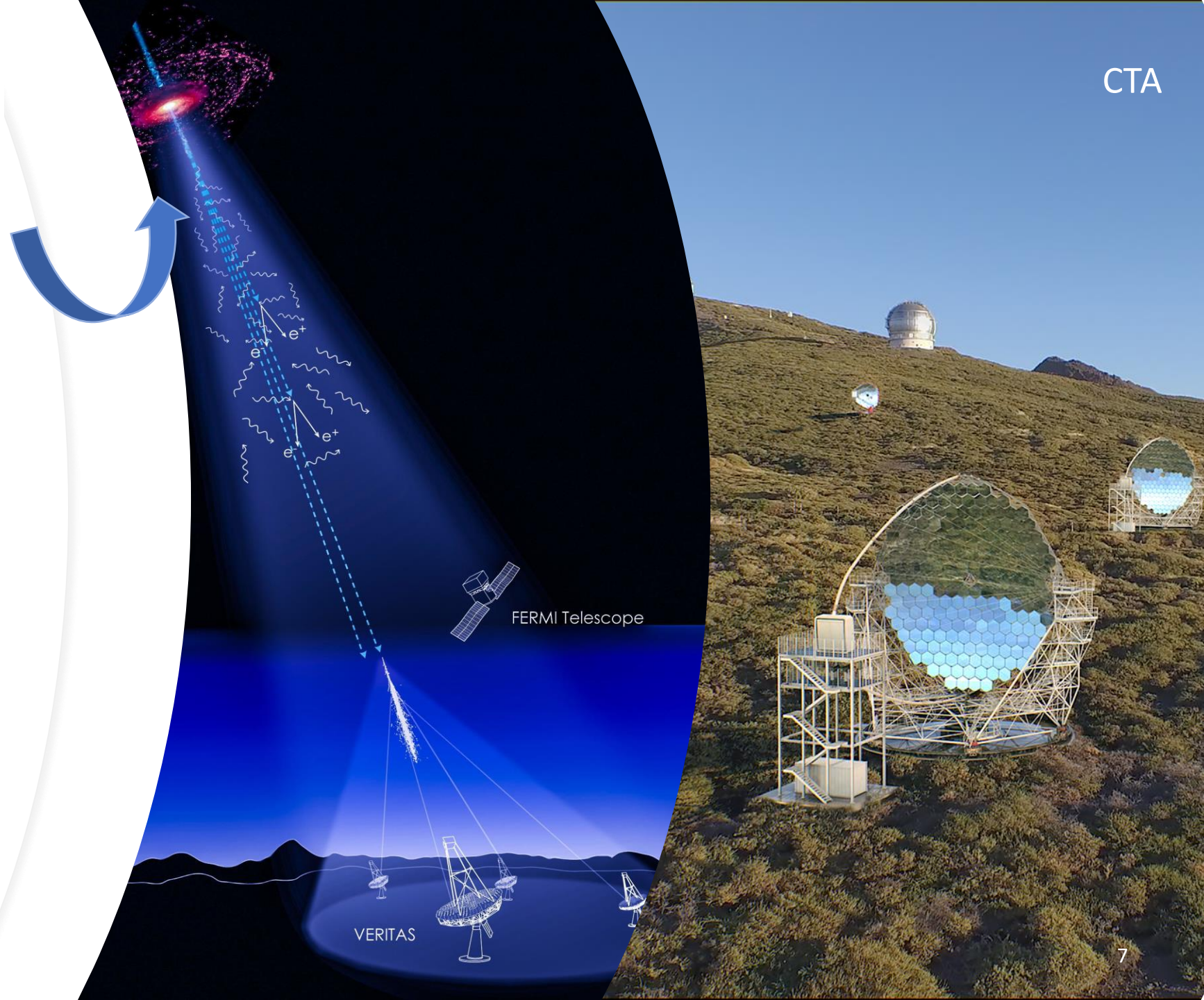


# Search for dark matter

- **Cosmological  $\Lambda$ CDM model**
- **Weakly Interacting Massive Particles (WIMPs)**
- **Particles are stable, massive, and interacts only through weak and gravitational interactions.**
- **Different approaches: Collider search, Direct detection and Indirect detection.**
- **Indirect: WIMPs annihilate into SM particles (bosons, quarks, leptons)**



# Detection of $\gamma$ -ray flux



# Expected $\gamma$ -ray flux

$$\frac{d\phi_\gamma}{dE_\gamma} = \frac{1}{\xi} \frac{\langle\sigma v\rangle}{4\pi m_\chi^2} \sum_f B_f \frac{dN_\gamma^f}{dE_\gamma} \times J$$

- $1/\xi$  depends on whether the DM particle is Majorana fermion or Dirac fermion
  - DM annihilation cross-section averaged over the velocity distribution  $\langle\sigma v\rangle$ 
    - DM particle mass  $m_\chi$
- J-factor describes the DM distribution and the amount of DM annihilations within the source



# Statistical analysis

- A widely used procedure to establish discovery is based on a frequentist significance test using a likelihood ratio as a test statistic.
- We use log-likelihood ratio statistical test to constrain the DM annihilation cross-section setting upper limits.

# Formalism

- Suppose, we are measuring  $x$ , and compute a histogram for all the events in the signal sample.
- The expectation value of number of events in each bin will be:

$$E[n_i] = \mu s_i + b_i$$

# Formalism

- Suppose, we are measuring  $x$ , and compute a histogram for all the events in the signal sample.
- The expectation value of number of events in each bin will be:

$$E[n_i] = \mu s_i + b_i$$

- The mean number of entries in the  $i_{th}$  bin from signal and background are:

$$s_i = s_{tot} \int_{bin\ i} f_s(x; \boldsymbol{\theta}_s) dx$$

$$b_i = b_{tot} \int_{bin\ i} f_b(x; \boldsymbol{\theta}_b) dx$$

# Formalism

- Now we compute the likelihood function which is the product of Poisson probabilities for all bins.

$$L(\mu, \boldsymbol{\theta}) = \prod_{j=1}^N \frac{(\mu s_j + b_j)^{n_j}}{n_j!} e^{-(\mu s_j + b_j)}$$

# Formalism


- Now we compute the likelihood function which is the product of Poisson probabilities for all bins.

$$L(\mu, \boldsymbol{\theta}) = \prod_{j=1}^N \frac{(\mu s_j + b_j)^{n_j}}{n_j!} e^{-(\mu s_j + b_j)}$$

- Further measurements to constrain the nuisance parameters:

$$L(\mu, \boldsymbol{\theta}) = \prod_{j=1}^N \frac{(\mu s_j + b_j)^{n_j}}{n_j!} e^{-(\mu s_j + b_j)} \prod_{k=1}^M \frac{u_k^{m_k}}{m_k!} e^{-u_k}$$

Background



# Formalism


Profile likelihood ratio

$$\lambda(\mu) = \frac{L(\mu, \hat{\hat{\theta}})}{L(\hat{\mu}, \hat{\theta})}$$

# Formalism

Profile likelihood ratio

$$\lambda(\mu) = \frac{L(\mu, \hat{\hat{\theta}})}{L(\hat{\mu}, \hat{\theta})}$$



Denotes the value of  $\theta$  that maximizes likelihood for the specified  $\mu$  (conditional)

# Formalism

Profile likelihood ratio

$$\lambda(\mu) = \frac{L(\mu, \hat{\hat{\theta}})}{L(\hat{\mu}, \hat{\theta})}$$

Denotes the value of  $\theta$  that maximizes likelihood for the specified  $\mu$  (conditional)

Unconditional ML estimators



# Formalism

- [G. Casella, R.L. Berger, *Statistical Inference*, Duxbury Press, 1990] -  $-2\ln\lambda$  converges in distribution to a  $\chi^2$  random variable with  $k$  degrees of freedom.

# Formalism

- [G. Casella, R.L. Berger, *Statistical Inference*, Duxbury Press, 1990] -  $-2\ln\lambda$  converges in distribution to a  $\chi^2$  random variable with  $k$  degrees of freedom.

$$t_{\mu} = -2 \ln \lambda(\mu)$$

$$p_{\mu} = \int_{t_{\mu, \text{obs}}}^{\infty} f(t_{\mu} | \mu) dt_{\mu}$$

# Formalism

$$t_{\mu} = -2 \ln \lambda(\mu)$$

- No condition on values of  $\mu$ .
- As  $\mu = 0$  : background-only events, presence of a new signal can only increase the mean event rate beyond what is expected from background alone.

# Formalism

$$t_{\mu} = -2 \ln \lambda(\mu)$$

- No condition on values of  $\mu$ .
- As  $\mu = 0$  : background-only events, presence of a new signal can only increase the mean event rate beyond what is expected from background alone.
- **$\mu \geq 0$**

# Formalism

$$\tilde{\lambda}(\mu) = \begin{cases} \frac{L(\mu, \hat{\boldsymbol{\theta}}(\mu))}{L(\hat{\mu}, \hat{\boldsymbol{\theta}})} & \hat{\mu} \geq 0 \\ \frac{L(\mu, \hat{\boldsymbol{\theta}}(\mu))}{L(0, \hat{\boldsymbol{\theta}}(0))} & \hat{\mu} < 0 \end{cases}$$

- For  $\hat{\mu} < 0$ , the best level of agreement between the data and any physical value of  $\mu$  occurs for  $\mu = 0$ .

# Formalism

$$\tilde{\lambda}(\mu) = \begin{cases} \frac{L(\mu, \hat{\boldsymbol{\theta}}(\mu))}{L(\hat{\mu}, \hat{\boldsymbol{\theta}})} & \hat{\mu} \geq 0 \\ \frac{L(\mu, \hat{\boldsymbol{\theta}}(\mu))}{L(0, \hat{\boldsymbol{\theta}}(0))} & \hat{\mu} < 0 \end{cases}$$

- For  $\hat{\mu} < 0$ , the best level of agreement between the data and any physical value of  $\mu$  occurs for  $\mu = 0$ .

$$\tilde{t}_{\mu} = -2 \ln \tilde{\lambda}(\mu) = \begin{cases} -2 \ln \frac{L(\mu, \hat{\boldsymbol{\theta}}(\mu))}{L(\hat{\mu}, \hat{\boldsymbol{\theta}})} & \hat{\mu} \geq 0 \\ -2 \ln \frac{L(\mu, \hat{\boldsymbol{\theta}}(\mu))}{L(0, \hat{\boldsymbol{\theta}}(0))} & \hat{\mu} < 0 \end{cases}$$

# Formalism : Upper limits

- To establish an upper limit on a parameter, let us define:

$$q_{\mu} = \begin{cases} -2 \ln \lambda(\mu) & \hat{\mu} \leq \mu \\ 0 & \hat{\mu} > \mu \end{cases}$$

# Formalism : Upper limits

- To establish an upper limit on a parameter, let us define:

$$q_{\mu} = \begin{cases} -2 \ln \lambda(\mu) & \hat{\mu} \leq \mu \\ 0 & \hat{\mu} > \mu \end{cases}$$

- To set an upper limit, we are only interested in the data corresponding to  $\mu >$  the ML estimator of  $\mu$ .



# Formalism : Upper limits

- But for the condition:  $\mu \geq 0$

$$q_\mu = \begin{cases} -2 \ln \lambda(\mu) \\ 0 \end{cases} \quad \hat{\mu} \leq \mu \quad = \quad \begin{cases} -2 \ln \frac{L(\mu, \hat{\boldsymbol{\theta}}(\mu))}{L(0, \hat{\boldsymbol{\theta}}(0))} & \hat{\mu} < 0 \\ -2 \ln \frac{L(\mu, \hat{\boldsymbol{\theta}}(\mu))}{L(\hat{\mu}, \hat{\boldsymbol{\theta}})} & 0 \leq \hat{\mu} < \mu \\ 0 & \hat{\mu} > \mu \end{cases}$$

$$p_\mu = \int_{q_{\mu, obs}}^{\infty} f(q_\mu | \mu) dq_\mu$$

# Examples: [C. Armand and B. Herrmann]

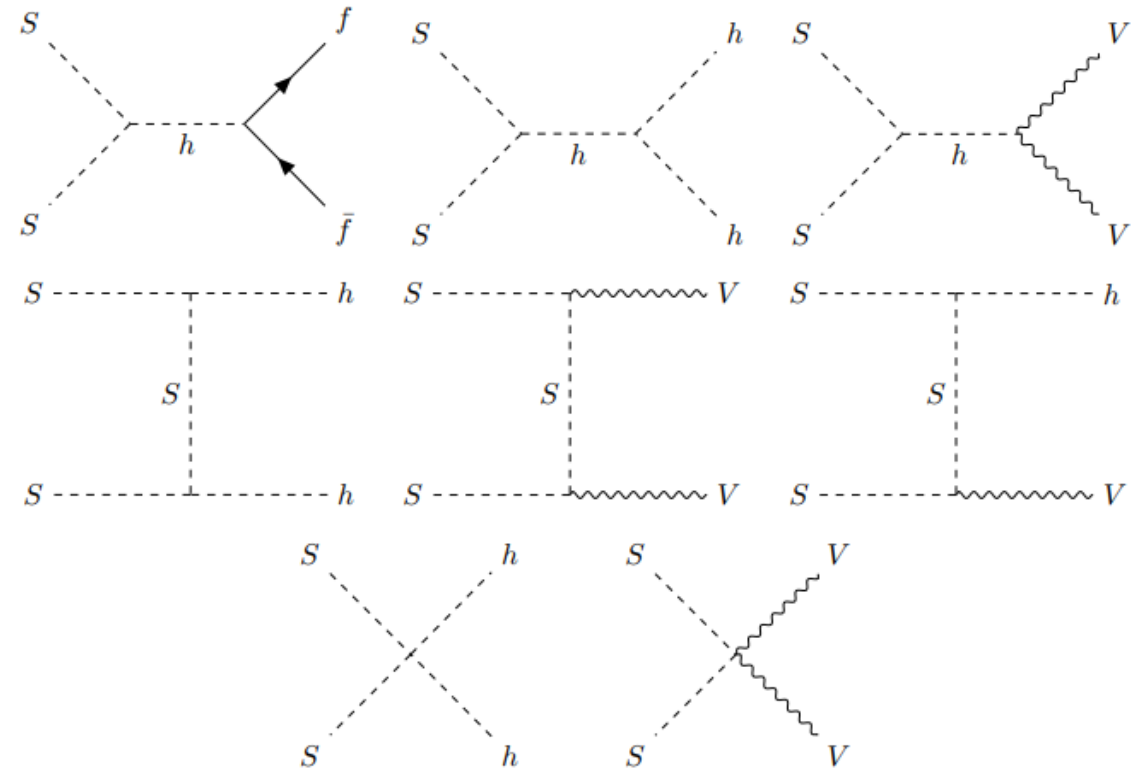
- 'Dark matter indirect detection limits from complete annihilation patterns'
- AIM: Studying simulated data of a dwarf galaxy to derive constraints on the DM annihilation cross-section within the singlet scalar dark matter model.
- Assumptions:
  - Candidate: real singlet scalar  $S$
  - DM annihilation into fermions proceeds solely through s-channel Higgs exchange.
  - DM annihilation into bosonic states can proceed through direct four-vertex interactions.

# Expected $\gamma$ -ray flux

$$\frac{d\phi_\gamma}{dE_\gamma} = \frac{1}{\xi} \frac{\langle\sigma v\rangle}{4\pi m_\chi^2} \sum_f B_f \frac{dN_\gamma^f}{dE_\gamma} \times J$$

- $1/\xi$  depends on whether the DM particle is Majorana fermion or Dirac fermion
  - DM annihilation cross-section averaged over the velocity distribution  $\langle\sigma v\rangle$ 
    - DM particle mass  $m_\chi$
- J-factor describes the DM distribution and the amount of DM annihilations within the source

# Feynman diagrams for all relevant annihilation products

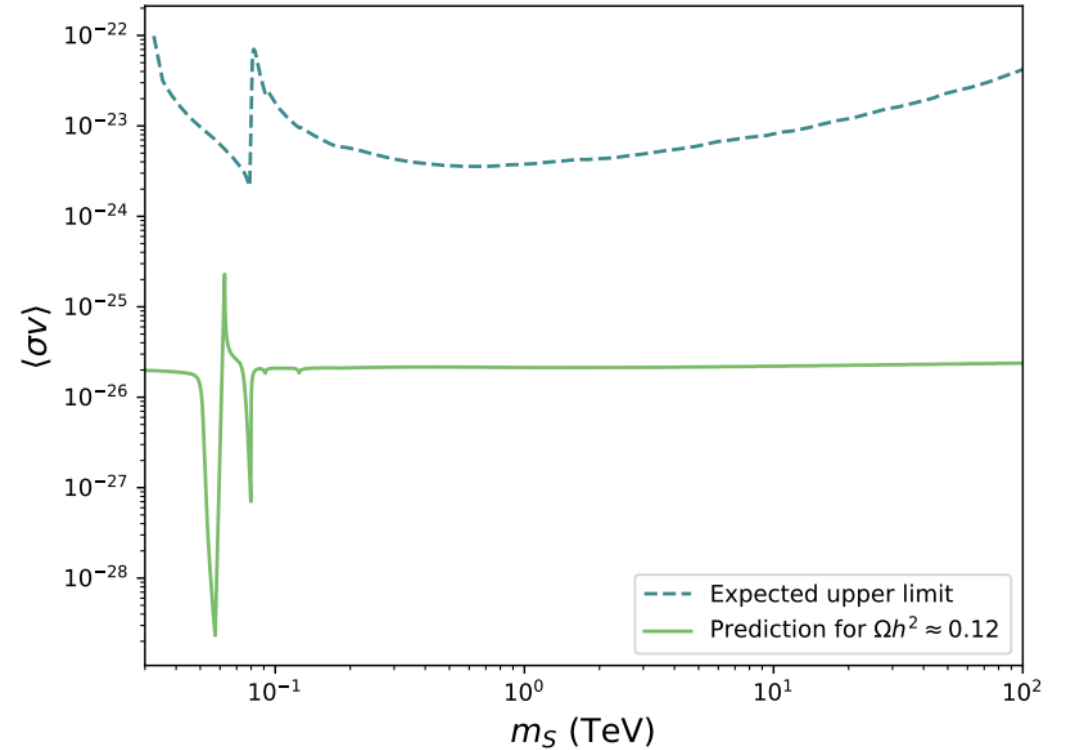
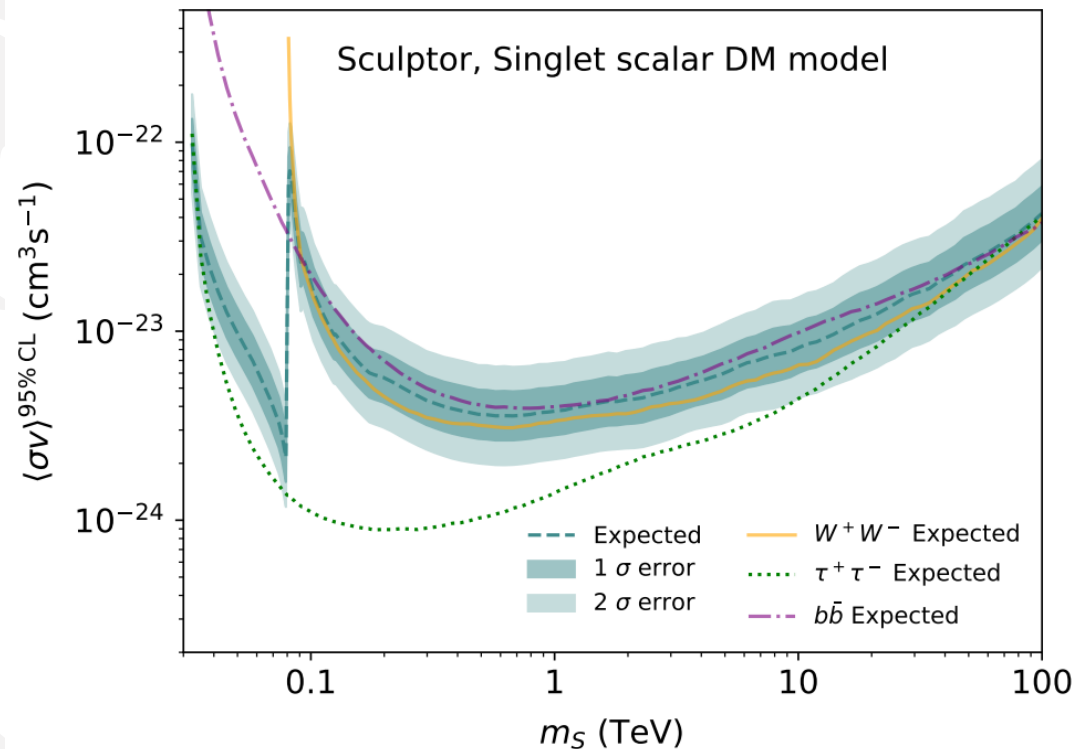


**Figure 1.** Tree-level Feynman diagrams for scalar pair annihilation into fermions ( $f = e, \mu, \tau, u, d, c, s, t, b$ ), gauge ( $V = W^\pm, Z^0$ ) and Higgs bosons ( $h$ ). Diagrams corresponding to  $u$ -channels obtained through crossing are not separately depicted. Annihilation into  $\gamma\gamma$  and  $gg$  final states proceeds through effective couplings to the Higgs boson  $h^0$ .

# Examples:

$$TS = \begin{cases} 0 & \text{for } \widehat{\langle \sigma v \rangle} > \langle \sigma v \rangle \\ -2 \ln \frac{\mathcal{L}(\langle \sigma v \rangle, \widehat{N}_B(\widehat{\langle \sigma v \rangle}), \widehat{J}(\widehat{\langle \sigma v \rangle}))}{\mathcal{L}(\widehat{\langle \sigma v \rangle}, \widehat{N}_B, \widehat{J})} & \text{for } 0 \leq \widehat{\langle \sigma v \rangle} \leq \langle \sigma v \rangle \\ -2 \ln \frac{\mathcal{L}(\langle \sigma v \rangle, \widehat{N}_B(\widehat{\langle \sigma v \rangle}), \widehat{J}(\widehat{\langle \sigma v \rangle}))}{\mathcal{L}(0, \widehat{N}_B(0), \widehat{J}(0))} & \text{for } \widehat{\langle \sigma v \rangle} < 0 \end{cases}$$

# Examples:

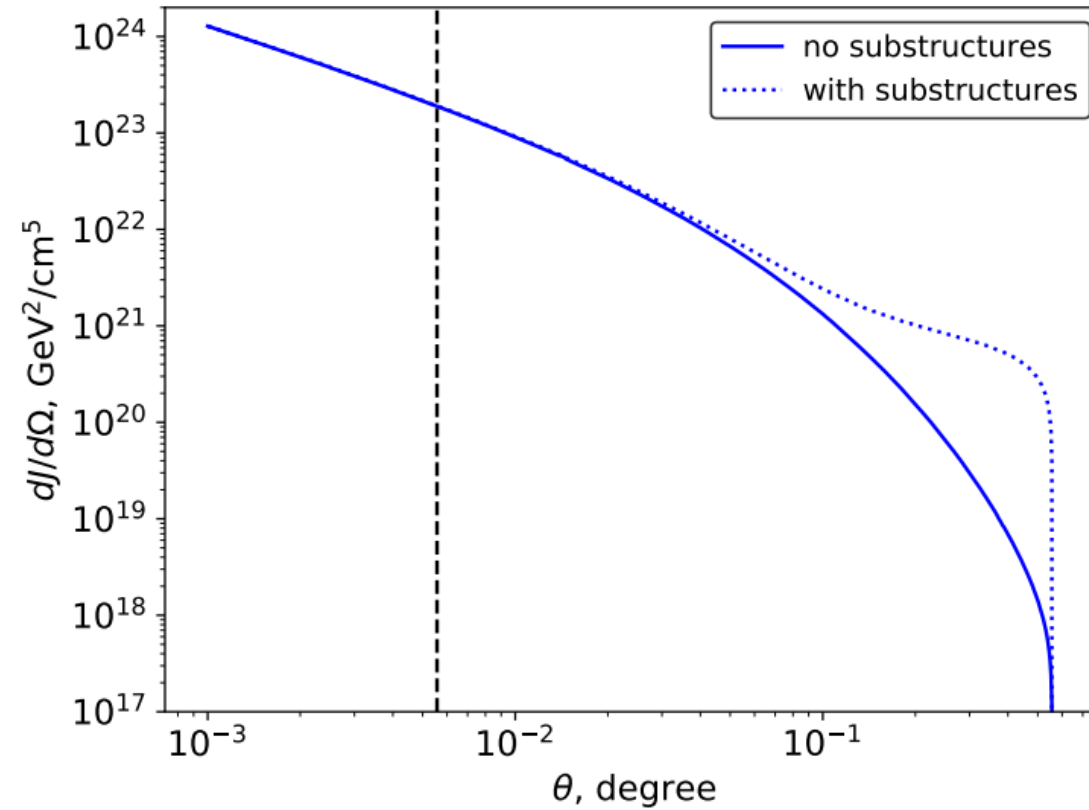


- $m_S = m_h/2 \approx 62.5 \text{ GeV}$  corresponds to Higgs resonance.
- $m_S \approx m_W \approx 80 \text{ GeV}$ , DM particles get heavier and produce more energetic SM particles

# Examples: [[Thorpe-Morgan et al.](#)]

- 'Annihilating Dark Matter Search with 12 Years of Fermi-LAT Data in Nearby Galaxy Clusters'
- AIM: Studying observed data of fermi-LAT of galaxy clusters to derive constraints on the DM annihilation cross-section to discuss the potential of a boost to the signal due to the presence of substructures in the DM halos
- Assumptions:
  - Cases with assuming the presence of substructures that can significantly boost the expected signal from the outer parts of halos.
  - The spatial and mass distributions are taken from [Sánchez-Conde & Prada \(2014\)](#)

# Examples:



$$dJ/d\Omega = \int_{l.o.s} \rho^2(l) dl$$



# Examples:

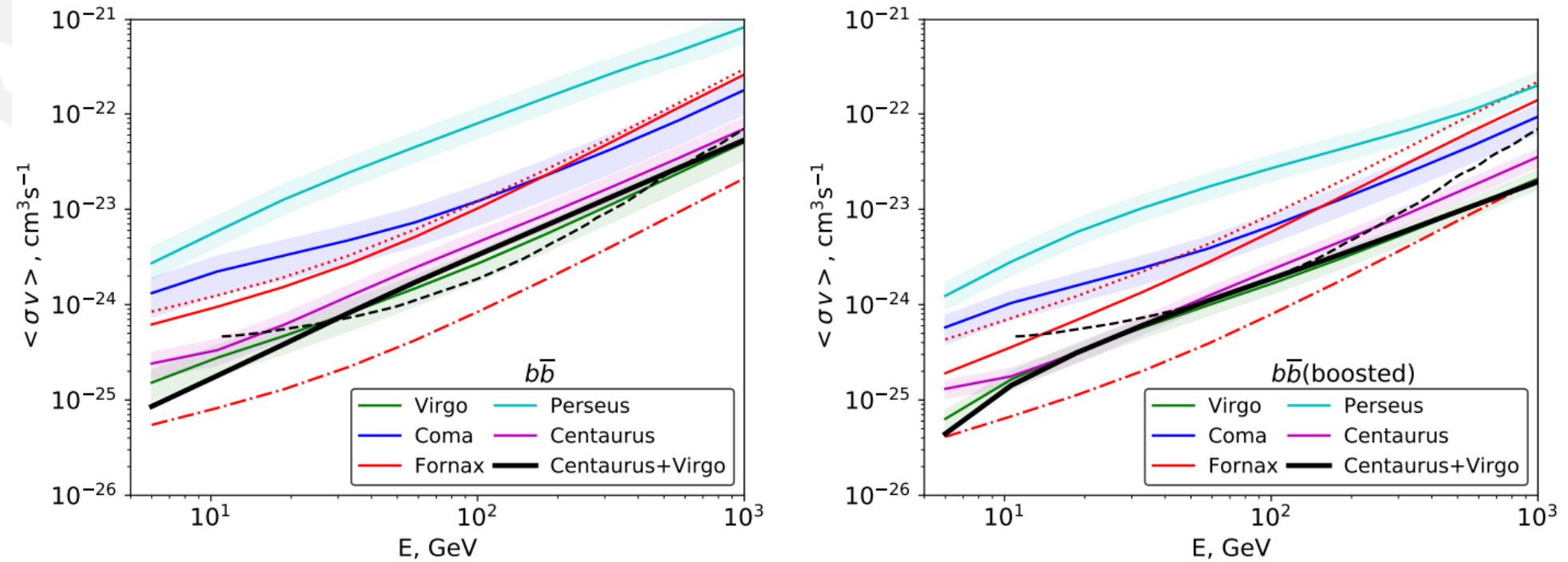


Fig: The left panel shows the results for smooth DM halos, while the right one shows the results when the presence of substructures is included

# Examples:

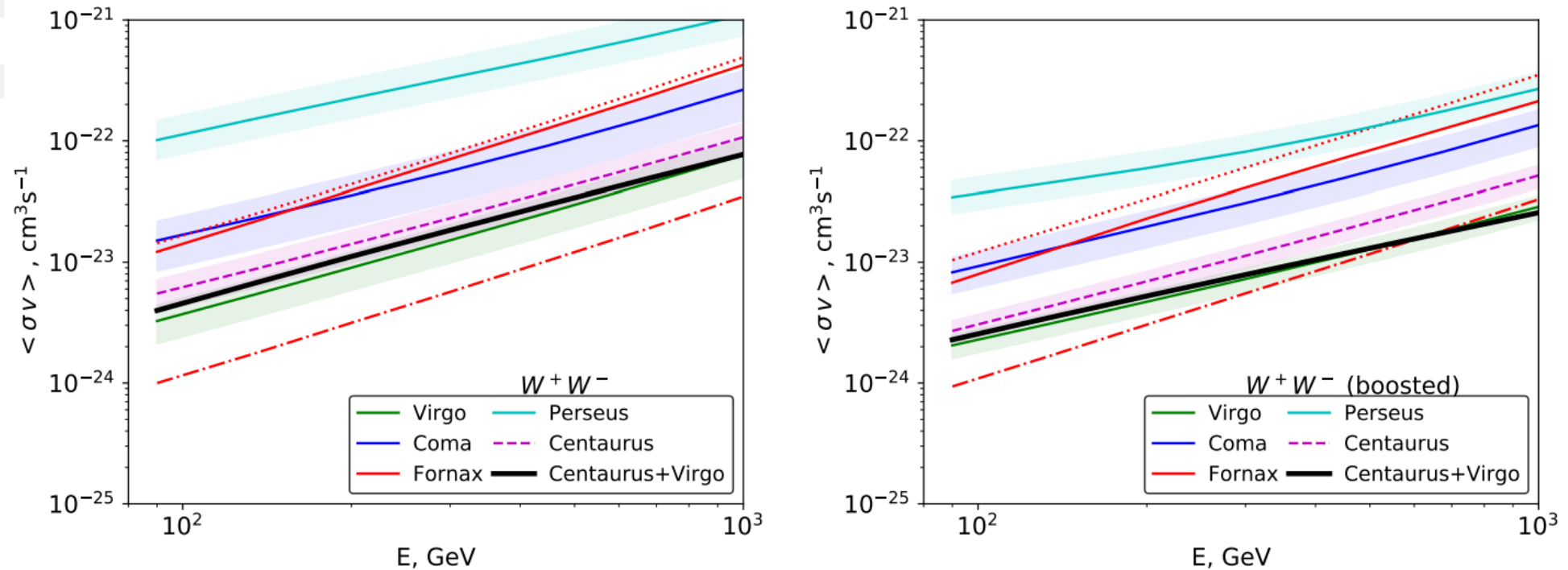


Fig: The left panel shows the results for smooth DM halos, while the right one shows the results when the presence of substructures is included

# References:

- [1] Asymptotic formulae for likelihood-based tests of new physics [\[url\]](#)
- [2] Dark matter indirect detection limits from complete annihilation patterns [\[url\]](#)
- [3] Annihilating Dark Matter Search with 12 Years of Fermi LAT Data in Nearby Galaxy Clusters [\[url\]](#)
- [4] Fundamental Physics Searches with IACTs [\[url\]](#)
- [5] Limits and confidence intervals in the presence of nuisance parameters [[W.A. Rolke et al. \(2005\)](#)]



Thank you

# Extra Slides

# Applications

- In this publication of 'Dark matter indirect detection limits from complete annihilation patterns':

$$\text{TS} = \begin{cases} 0 & \text{for } \widehat{\langle\sigma v\rangle} > \langle\sigma v\rangle, \\ -2 \ln \frac{\mathcal{L}(\langle\sigma v\rangle, \hat{\mathbf{N}}_{\text{B}}(\langle\sigma v\rangle), \hat{\mathbf{J}}(\langle\sigma v\rangle))}{\mathcal{L}(\widehat{\langle\sigma v\rangle}, \hat{\mathbf{N}}_{\text{B}}, \hat{\mathbf{J}})} & \text{for } 0 \leq \widehat{\langle\sigma v\rangle} \leq \langle\sigma v\rangle, \\ -2 \ln \frac{\mathcal{L}(\langle\sigma v\rangle, \hat{\mathbf{N}}_{\text{B}}(\langle\sigma v\rangle), \hat{\mathbf{J}}(\langle\sigma v\rangle))}{\mathcal{L}(0, \hat{\mathbf{N}}_{\text{B}}(0), \hat{\mathbf{J}}(0))} & \text{for } \widehat{\langle\sigma v\rangle} < 0, \end{cases}$$

# Formalism

$$\lambda(\mu) = \frac{L(\mu, \hat{\theta})}{L(\hat{\mu}, \hat{\theta})}$$

- Here,  $\hat{\theta}$  denotes the value of  $\theta$  that maximizes  $L$  for the specified  $\mu$ . (conditional maximum-likelihood (ML) estimator)
- The denominator is the maximized (unconditional) likelihood function.
- We generally define the test statistic as:

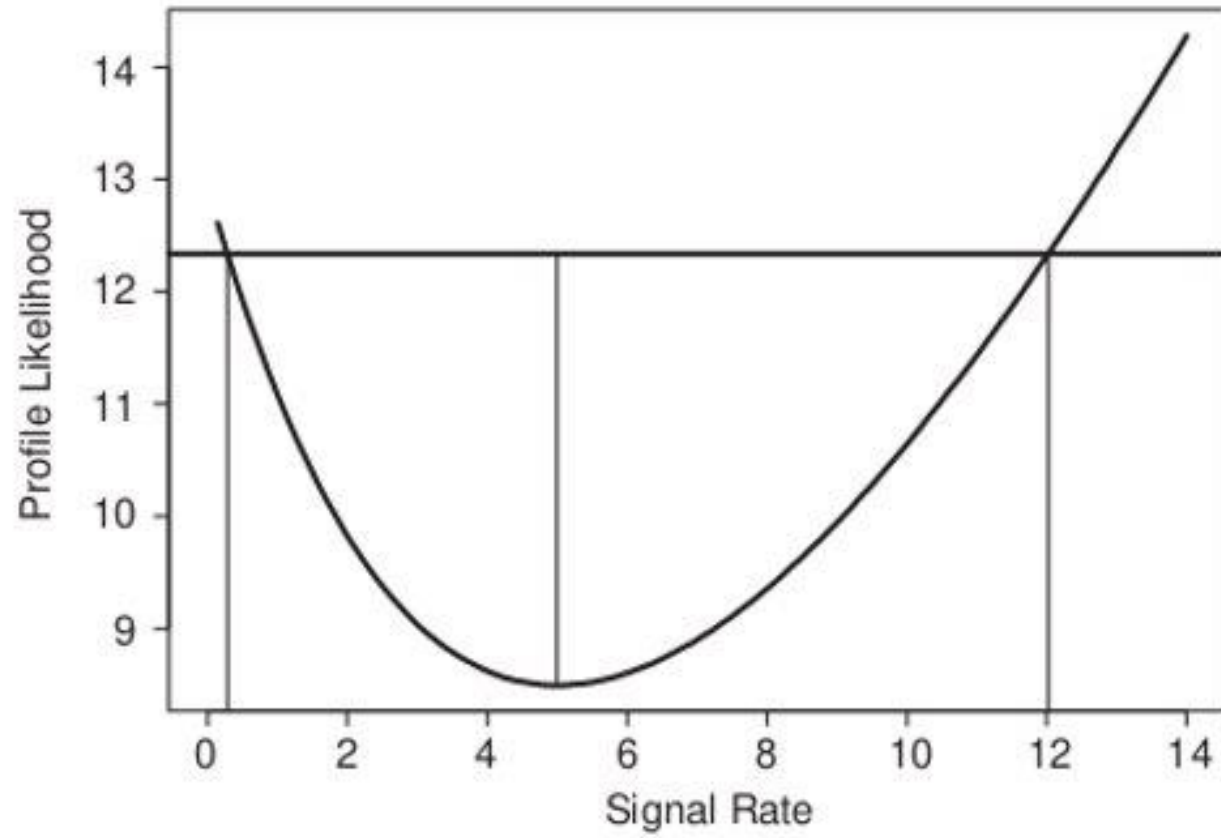
$$t_{\mu} = -2 \ln \lambda(\mu)$$

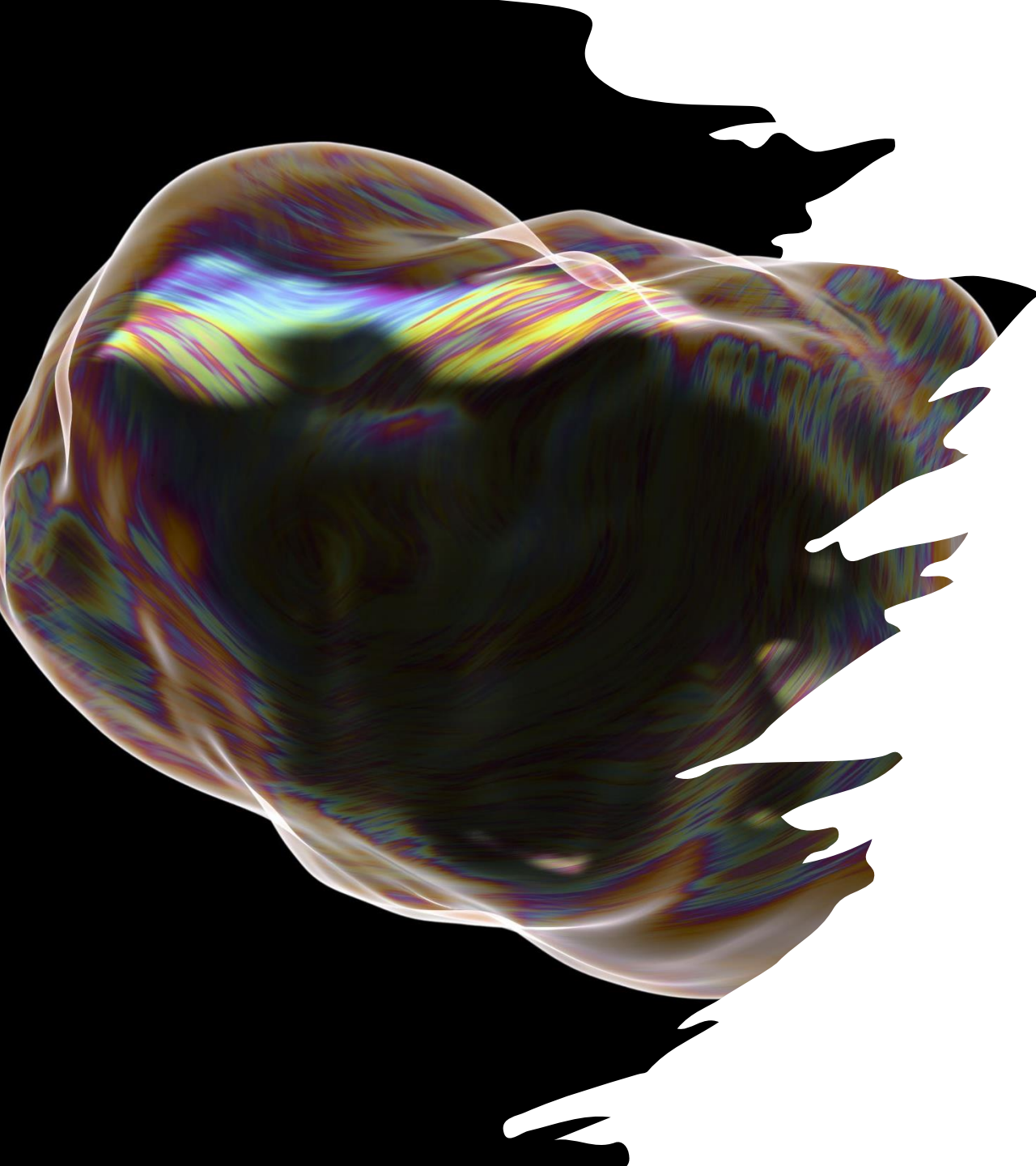
As when using  $q_\mu$ , the upper limit on  $\mu$  at confidence level  $1 - \alpha$  is found by setting  $p_\mu = \alpha$  and solving for  $\mu$ , which reduces to the same result as found when using  $q_\mu$ , namely,

$$\mu_{\text{up}} = \hat{\mu} + \sigma \Phi^{-1}(1 - \alpha) . \quad (69)$$

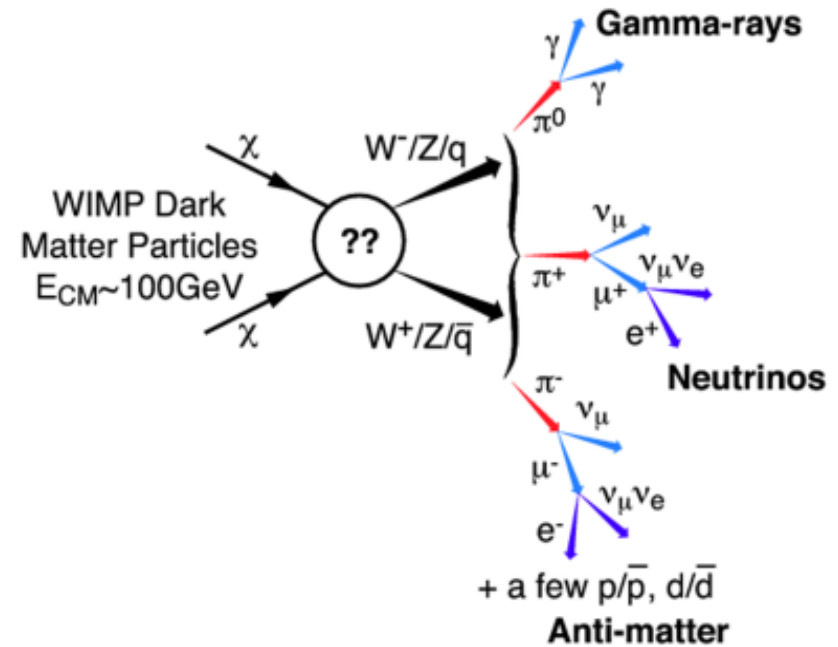
That is, to the extent that the Wald approximation holds, the two statistics  $q_\mu$  and  $\tilde{q}_\mu$  lead to identical upper limits.







# Search for dark matter



- A widely used procedure to establish discovery is based on a frequentist significance test using a likelihood ratio as a test statistic.
- We use log-likelihood ratio statistical test to constrain the DM annihilation cross-section setting upper limits.

## Statistical analysis