Derivation of spectral cutoff lower limits in PeVatron searches with CTA

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1. Introduction

- PeVatron is a term used to describe astrophysical sources that are able to accelerate particles up to $10^{15} eV$ (1 PeV)
- Several source classes have been proposed as potential PeVatrons, but Supernova Remnants (SNRs) have been the preferred candidates
- Galactic PeVatron have been detected, but none of them are proven to be related to SNRs
- Crab nebula is an example of a leptonic PeVatron (in this work hadronic PeVatrons are being searched for)

By NASA, ESA, J. Hester and A. Loll (Arizona State University)

- CTA is the next generation Imaging Atmosferic Cherenkov Telescope (IACT) system.
- It will be located at Paranal Observatory (Chile) and Roque de los Mucachos Observatory (Spain)→whole sky observations
- Energy range from 20 GeV to 200 TeV
- Improved sensitivity is expected to lead to discovery of many more astrophysical sources.

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2. Simulation and analysis of CTA data

- The simulation and analysis of CTA dana are based on the instrument response functions (IRFs)
- The morphology of extended γ-ray sources is modelled using 2D symmetric Gaussians, and source extensions are given as the width (σ) of the Gaussian.
- Simulated CTA event data are drawn from Poisson distributed random variables around their bin-wise expectation.
- A binned 3D-likelihood analysis is performed in the framework of gammapy.
- The population of Galactic SNRs is simulated with a Monte Carlo approach, in which the distribution of SNe in time and space is randomly drawn in multiple samples.

3. Derivation of spectral cutoff lower limits

- PeVatron searches with CTA rely on the derivation of statistical statements on the inverse energy cutoff parameter λ .
- When a significant cutoff detection is impossible, frequentist upper limits λ^{UL} on the inverse cutoff parameter at a given confidence level CL are of high relevance.
- Limits on the inverse spectral cutoff λ are investigated within γ -ray emission models.

The Poisson likelihood:

Cash statistic:

$$L(\lambda, \theta | \vec{c}) \coloneqq \prod_{i=1}^{N} \exp(-n_i) \frac{n_i^{c_i}}{c_i!}$$

$$C(\lambda,\theta) = 2\sum_{i} (n_i - c_i \ln n_i)$$

- θ nuisance parameters
- λ the inverse energy cutoff
- $\vec{c} = (c_1, ..., c_N)$ simulated event counts
- $\vec{n} = \vec{n}(\lambda, \theta)$ predicted counts

3.1. Profile likelihood

 This method is an example for the inversion of a frequentist hypothesis test

The likelihood ratio test statistic: $\Lambda(\lambda) \coloneqq -2 \ln \frac{L(\lambda)}{L(\hat{\lambda})} = C(\lambda) - C(\hat{\lambda})$

• $\hat{\lambda}$ - maximum likelihood estimator for the inverse energy cutoff over the constrained range $\lambda \ge 0$

• Let $L(\lambda) = max_{\theta}L(\lambda, \theta | \vec{c})$ be profile likelihood, and $C(\lambda) = max_{\theta}C(\lambda, \theta)$ corresponding Cash statistic Likelihood ratio statistic for the analysis of a typical γ -ray source:



- The likelihood ratio test statistic equation is the test statistic of the null hypothesis (H_0 : $\lambda = \hat{\lambda}$) against the alternative hypothesis (H_1 : $\lambda = \lambda$)
- The alternative hypothesis is accepted when the test statistic is smaller than or equal to the critical value at a given confidence level CL (the horizontal line in the previous slide)

3.2. Markov-Chain Monte Carlo (MCMC)

 Upper limit of the inverse cutoff parameter can be derived when the probability distribution of model parameters is expressed in the framework of Bayesian terminology.

Posterior probability density

$$p(\lambda, \theta | \vec{c}) = \frac{L(\lambda, \theta | \vec{c}) p(\lambda, \theta)}{p(\vec{c})}$$

- $p(\lambda, \theta)$ probability density for the model parameters
- $p(\vec{c}) \coloneqq \int d\lambda d\theta \ L(\lambda, \theta | \vec{c}) \ p(\lambda, \theta)$

3.3. Bootstrap

- This method resamples binned γ -ray events (\vec{c}) as bootstrap samples (\vec{c}_*)
- The percentile method is used to get the smallest positive upper limit on the inverse cutoff parameter (λ^{UL}) which satisfies:

$$CL \leq \int_{-\infty}^{\lambda^{UL}} d\lambda_* f(\lambda_*)$$

- Non-parametric bootstrap
- Parametric bootstraps:
 - Poisson bootstrap
 - Best fit bootstrap
- Difference between the parametric and non-parametric bootstrap is that the total number of events is a random variable

3.4. Performance comparison

Point-like source analysis

- The energy cutoff limits obtained with the bootstrap and MCMC methods are calculated with a precision better than 2%.
- Two different sets of prior density distributions for the model parameters are investigated for the MCMC method.
- The uniformity of prior density depends on the choice of the parameter.
- Results based on uniform prior densities are compared to results obtained with priors based on gamma distributed random variables.





A comparison of the coverage and sensitivity for the specific γ -ray point-like source



Point-like source energy cutoff sensitivity as a function of the true energy cutoff for different methods relative to the respective sensitivity achievable with a 1-dimensional profile likelihood analysis.

The different panels show the relative sensitivity for different point-like source parameters in terms of flux normalization ϕ_0 and index Γ .

Conclusions

- The profile likelihood method provides a computationally very efficient way to derive lower limits on the energy cutoff.
- Other methods are less sensitive or a possible sensitivity improvement in restricted parameter ranges results from the choice of the prior distributions(MCMC).
- The computational effort to derive reasonably precise limits is larger for bootstrap and MCMC implementations than for the profile likelihood method.
- Bootstrap and MCMC methods can provide an important alternative in cases where the profile likelihood method cannot be applied.

Bibliography

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Thank you for the attention